

Large Density Phase of FT Two-Color QCD

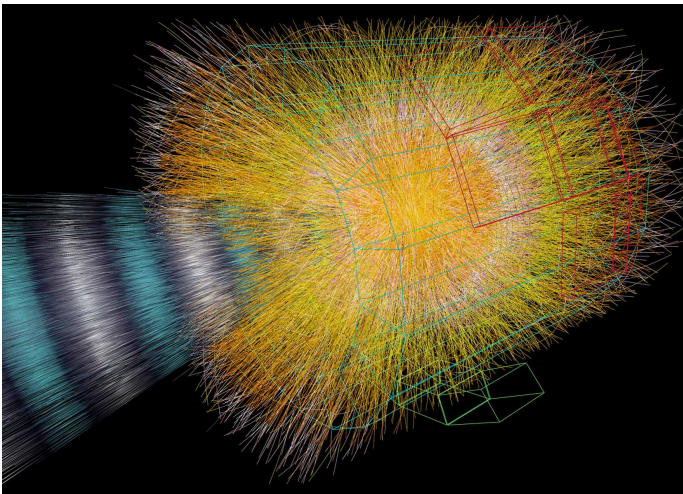
Seyong Kim

Sejong University

Outline

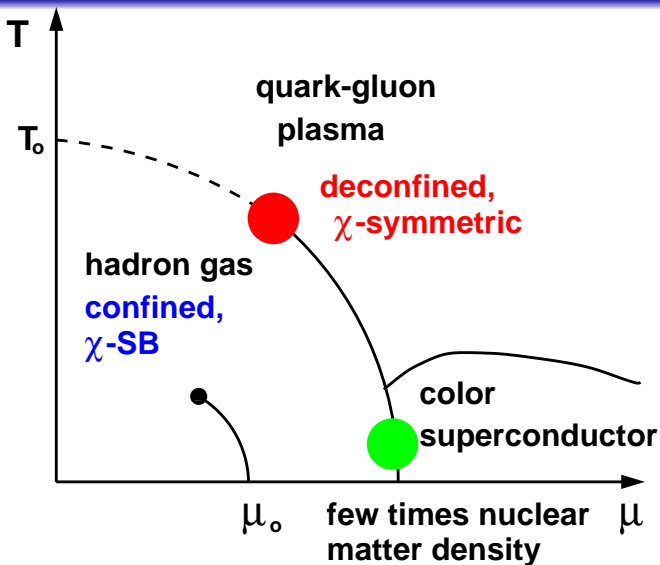
- 1 Intro
- 2 F.T./F.D.
- 3 SU(2)
- 4 low T, μ
- 5 finite T, μ
- 6 discussion

Motivation

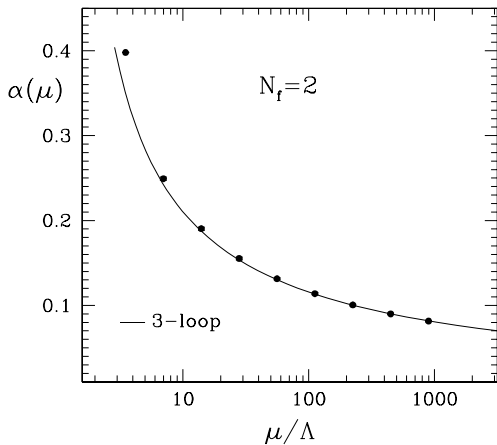


from <http://doc.cern.ch/archive/electronic/cern/others/PHO/photo-bul/bul-pho-2007-073.jpg>

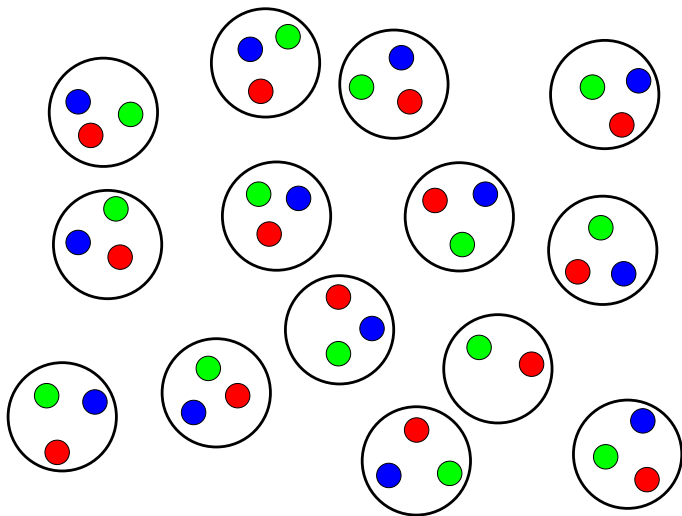
Motivation



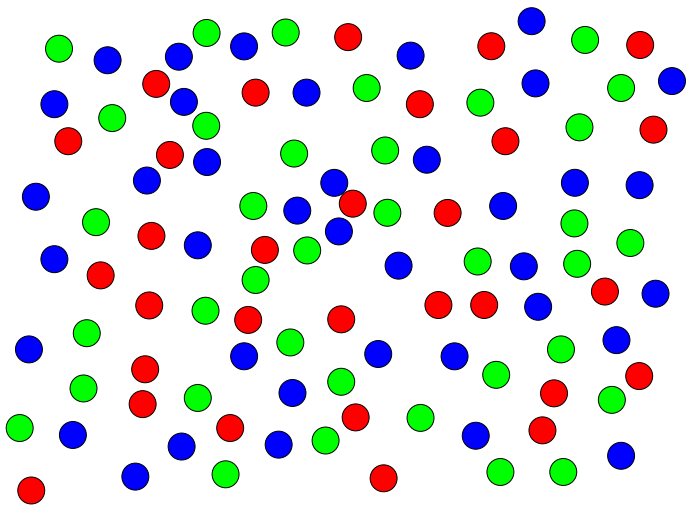
Motivation



Motivation



Motivation



Finite Temperature/Density Lattice Field Theory

- statistical mechanics
- consider Ising model

$$Z = \text{Tre}^{-H} \quad (1)$$

where

$$H = -J \sum_x \sigma_x \sigma_{x+i} - h \sum_x \sigma_x \quad (2)$$

with $\sigma_x = \pm 1$

Finite Temperature/Density Lattice Field Theory

- symmetry \rightarrow order parameter

$$\begin{aligned}\langle M \rangle &= \left\langle \sum_x \sigma_x \right\rangle \\ &= \frac{\text{Tr} \sum_x \sigma_x e^{-H}}{\text{Tr} e^{-H}}\end{aligned}\quad (1)$$

- non-symmetry related quantity

$$\begin{aligned}\langle E \rangle &= \langle H \rangle \\ &= \frac{\text{Tr} H e^{-H}}{\text{Tr} e^{-H}}\end{aligned}\quad (2)$$

Finite Temperature/Density Lattice Field Theory

- $k_B = 1$

$$\begin{aligned} Z_C &= \text{Tr} e^{-\frac{H}{T}} \\ &= \sum_{\phi} \langle \phi | e^{-\frac{H}{T}} | \phi \rangle \end{aligned} \tag{1}$$

Finite Temperature/Density Lattice Field Theory

- recall

$$F(q_f, q_i) = \langle q_f | e^{-iHt} | q_i \rangle \quad (1)$$

$$= \int dq_1 dq_2 \cdots dq_{N-1} \langle q_f | \hat{T} | q_{N-1} \rangle \langle q_{N-1} | \hat{T} | q_{N-2} \rangle \cdots \langle q_1 | \hat{T} | q_i \rangle$$

$$\rightarrow \int \mathcal{D}x(t) e^{i \int dt L}$$

$$\rightarrow \int \mathcal{D}x(\tau) e^{-\int d\tau L_E} \quad (2)$$

under $t \rightarrow i\tau$

Finite Temperature/Density Lattice Field Theory

- bosonic field: periodic in the time direction
- fermionic field: anti-periodic in the time direction

$$Z_C = \sum_{\phi} \langle \phi | e^{-S_E} | \phi \rangle \quad (1)$$

where

$$S_E = \int d^3x \int_0^{1/T} L_E \quad (2)$$

Finite Temperature/Density Lattice Field Theory

$$S_E = -\frac{1}{4} \text{Tr} \log(M^\dagger M) + \sum_{x, \mu\nu} \frac{1}{6g^2} \text{Tr}[1 - P_{\mu\nu}(x)] \quad (1)$$

where

$$M_{x,y} = \frac{1}{2} \sum_{x,\mu} \bar{\chi}(x) \eta_\mu(x) [U_\mu(x) \chi(x + \hat{\mu}) - U_\mu^\dagger(x - \hat{\mu}) \chi(x - \hat{\mu})] \quad (2)$$

$\eta_\mu(x)$ is the Kawamoto-Smit phase factor in the staggered fermion method

Finite Temperature/Density Lattice Field Theory

- symmetry \rightarrow order parameter
- chiral symmetry

$$\langle \bar{\psi}\psi \rangle \quad (1)$$

- Z(3) symmetry for quenched theory

$$\langle \prod_t U_t(x, t) \rangle \quad (2)$$

Finite Temperature/Density Lattice Field Theory

$$Z_{GC} = \text{Tre}^{-\frac{H-\mu N}{T}} \quad (1)$$

where $N = \bar{\Psi}\gamma_4\Psi$

- F. Karsch and P. Hasenfratz, PLB125, 308(1983)

$$M_{x,y} = \frac{1}{2} \sum_{x,\mu} \bar{\chi}(x) \eta_\mu(x) [U_\mu(x) e^{\mu} \chi(x + \hat{\mu}) - U_\mu^\dagger(x - \hat{\mu}) e^{-\mu} \chi(x - \hat{\mu})] \quad (2)$$

Finite Temperature/Density Lattice Field Theory

- complex action problem
 - model study

2-color gauge theory

- Two color QCD or SU(2) gauge theory in large chemical potential is different from QCD

- For SU(2),

$$(C\gamma_5)\tau_2 M(\mu)(C\gamma_5)^{-1}\tau_2 = M^*(\mu) \quad (1)$$

- $\det M(\mu)$ is real (but does not mean that it is positive)
- There is spontaneous chiral symmetry breaking
→ pion is light
- **But** qq is a color singlet → diquark condensate does not break color symmetry

2-color gauge theory

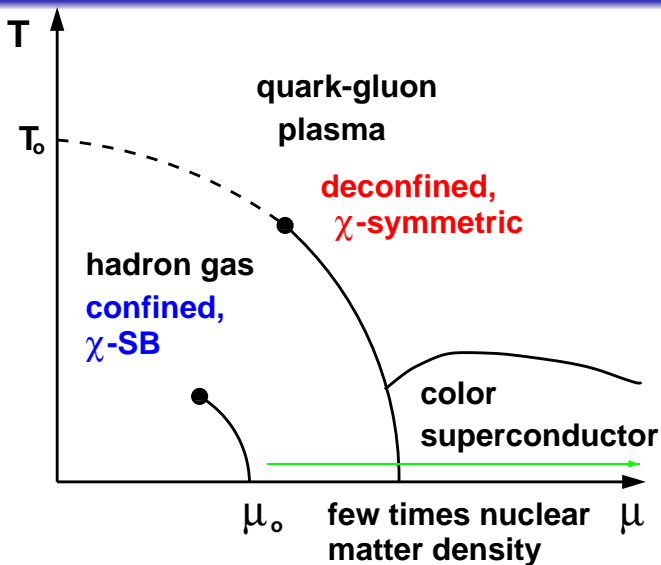
- Model for QCD but
- gluon sector is similar to QCD
- large chemical potential region can be studied

2-color gauge theory

- BEC phase : Kogut et al, Nucl. Phys. B582 (2000)477
- $n_q \propto f_\pi^2 (\mu - \mu_0)$
- $\langle qq \rangle \propto \sqrt{1 - \left(\frac{\mu_0}{\mu}\right)^4}$

2-color gauge theory

- BCS phase
- $n_q \propto \mu^3$
- $\varepsilon_F \propto \mu^4$
- $\langle qq \rangle \propto \Delta \mu^2$

Low T, μ 

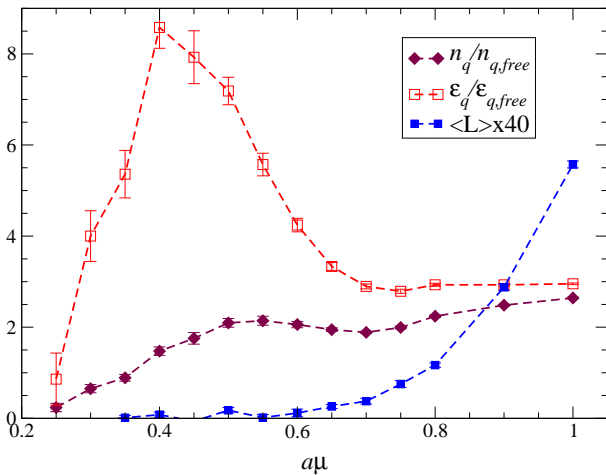
Low T, μ

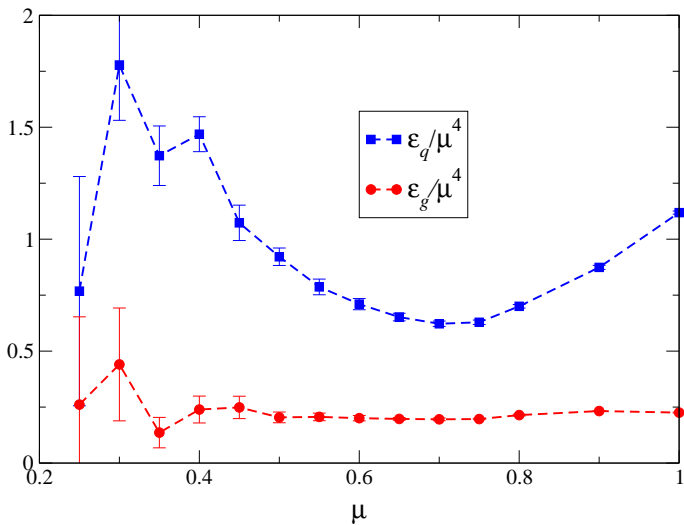
- Two color QCD with heavy quark:

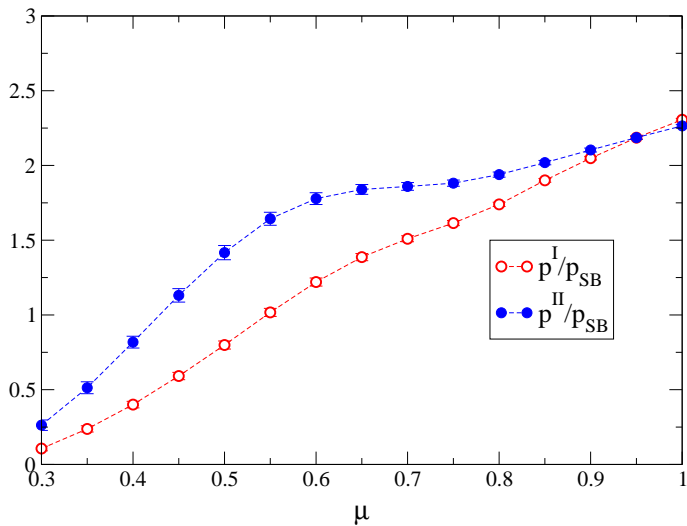
S. Hands, S. K., J.-I. Skullerud, Eur.Phys.J.C48:193,2006

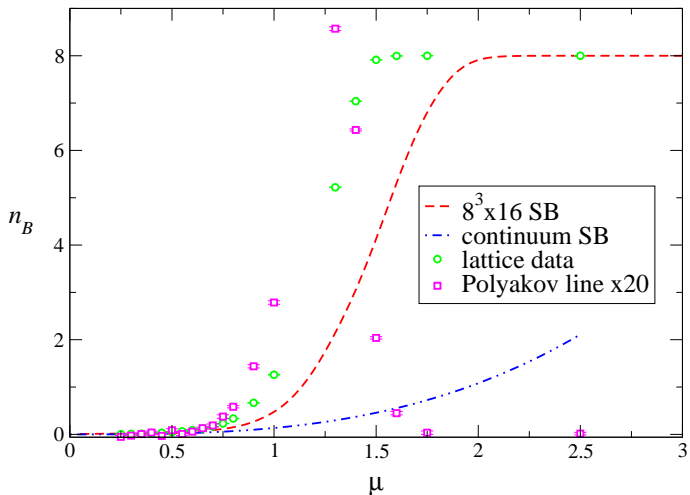
$8^3 \times 16, \beta = 1.7$ Wilson quark

$\kappa = 0.168, m_\pi a \sim 0.8$

Low T, μ 

Low T, μ 

Low T, μ 

Low T, μ 

Low T, μ

- Hint of three different phases

Hadronic phase

Bose-Einstein Condensed(BEC) phase

Bardeen-Cooper-Schrieffer(BCS) phase

- BEC-BCS transition is **not a sharp** transition

strong attraction \rightarrow tightly bound boson

\rightarrow boson condensation

weaker attraction \rightarrow loosely bound Cooper pair

\rightarrow superconducting phase

Low T, μ

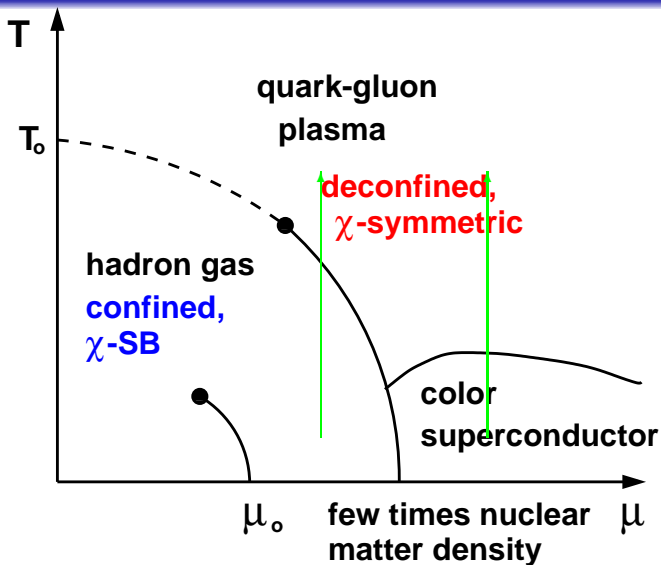
- P. Nozieres, S.Schmitt-Rink, J. of Low Temp. Phys. 59, 195(1985)

finite temperature transition of BEC phase is different

$$T_c = (2\pi/M)(N_p/2.612)^{2/3}$$

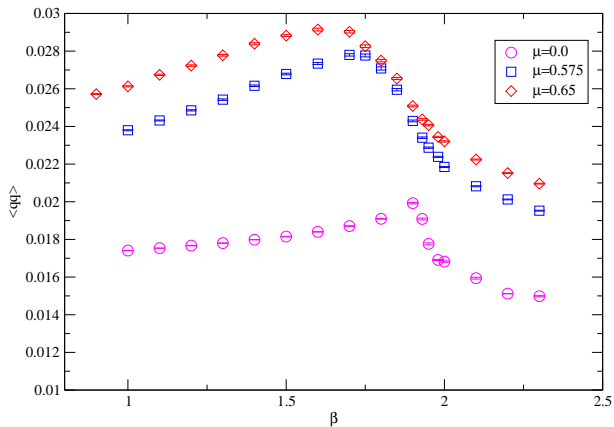
from that of BCS phase

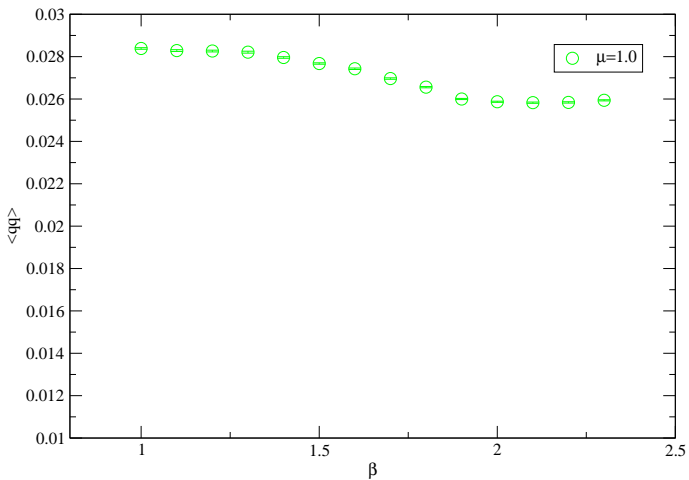
$$T_c = (e^\gamma/\pi)\Delta_F$$

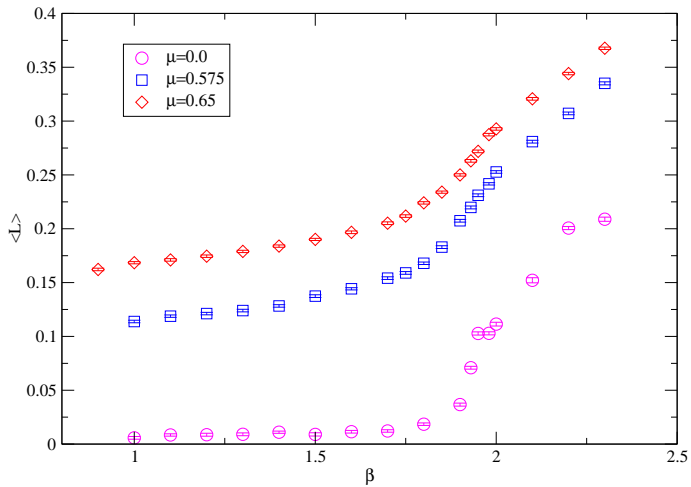
Finite T, μ 

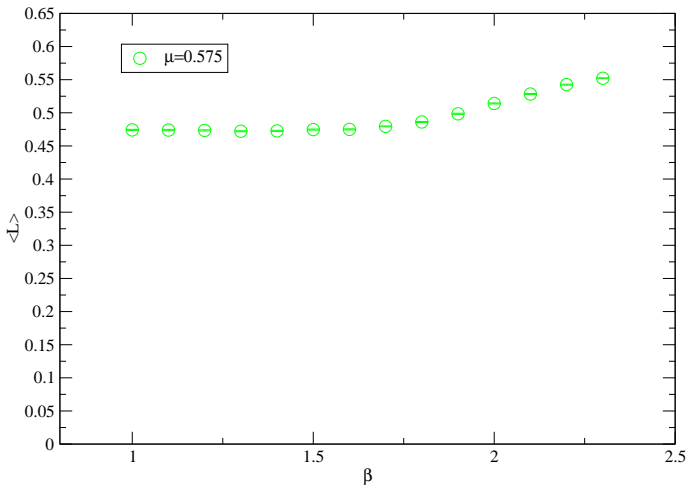
Finite T, μ

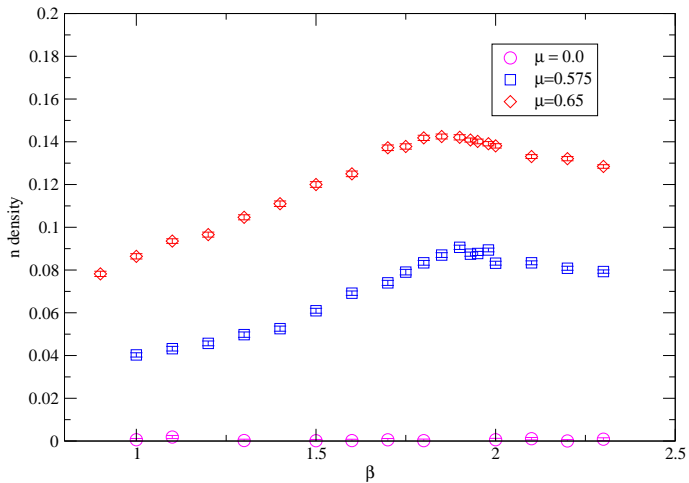
- temperature is defined as $T = \frac{1}{N_\tau a}$
 - ← a is controlled by gauge coupling constant
 - ← different temperature means a different gauge coupling constant
- periodic boundary condition on bosonic field, anti-periodic boundary condition on fermionic field
- no. of spatial lattice sites (N_S) should be bigger than N_τ
- phases are distinguished by order parameters

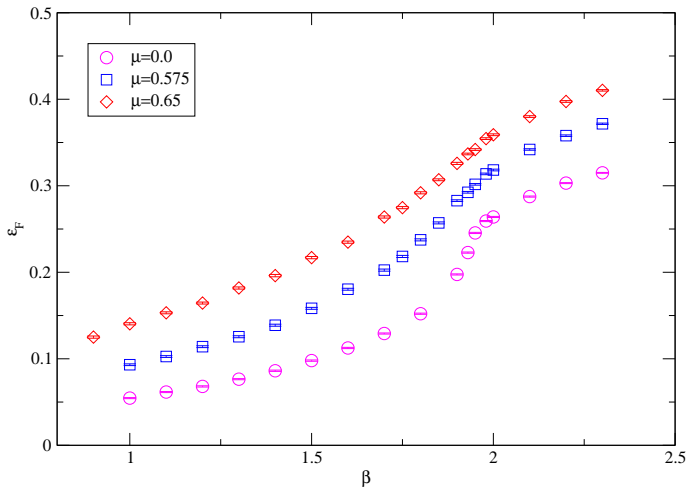
Finite T, μ 

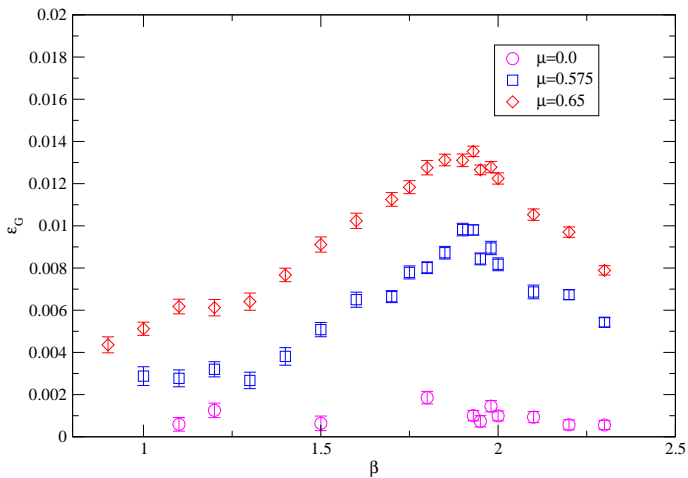
Finite T, μ 

Finite T, μ 

Finite T, μ 

Finite T, μ 

Finite T, μ 

Finite T, μ 

Discussion

- diquark condensate behavior shows interesting behavior

for $\mu = 0.575, 0.65$, the diquark condensate doesn't change much

for $1.75 < \beta < 2.0$

- further study is needed
- many parameters to scan \rightarrow Grid computing

Discussion

